

# LS depolarizing transitions in A+H collisions: a diabatization approach

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### **Motivation**



#### Discrepancy within calculations - Depolarization rates



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Why are they different?, Potential interaction

- Ab initio calculations Adiabatic representation
- Rayleigh-Schrödinger-Unsöld (RSU)



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#### Other sources

- Semiclassical theory
- Quantum dynamics theory

### Literature



#### Semiclassical - Pertubation theory potential interaction

Perturbation theory potentials - Single electron model

- Brueckner, Anstee & O'Mara, Derouich et al., ...
- Good agreement at intermediate-long distances
- Some anomalies could arise because of the effect of avoided crossings and ionic effects

### Literature



#### Semiclassical - Pertubation theory potential interaction

Perturbation theory potentials - Single electron model

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- Good agreement at intermediate-long distances
- Some anomalies could arise because of the effect of avoided crossings and ionic effects

#### Quantum - Adiabatic potential interaction

- Kerkeni and us
- Discrepancies due to curve crossing effects

### Avoided crossings





### Adiabatic curves





#### **Abnormalies**

- Avoided crossings
- M<sup>+</sup> + H<sup>-</sup> states cross with M(*L*, *S*) + H

Born-Oppenheimer approximation TI Schrödinger equation -  $H\Psi(r, R) = E\Psi(r, R)$  $H = -\frac{\hbar^2}{2} \sum_{i}^{n_N} \frac{1}{m_i} \nabla_{R_i}^2 - \frac{\hbar^2}{2m_e} \sum_{i}^{n_e} \nabla_{r_i}^2 + V(r_1, ..., r_{n_e}, R_1, ..., R_{n_N})$  $V = q_{e}^{2} \left| -\sum_{i}^{n_{N}} \sum_{i}^{n_{e}} \frac{\overline{\mathcal{Z}_{i}}}{|R_{i} - \underline{r_{j}}|} + \sum_{i}^{n_{N}} \frac{\overline{\mathcal{Z}_{i}}\overline{\mathcal{Z}_{j}}}{|R_{i} - R_{j}|} + \sum_{i}^{n_{e}} \frac{1}{|r_{i} - r_{j}|} \right|$ 

### Born-Oppenheimer approximation



TI Schrödinger equation -  $H\Psi(r, R) = E\Psi(r, R)$ 

$$egin{aligned} \mathcal{H}_{e} &= \, - \, rac{\hbar^2}{2 \, m_e} \sum_{i}^{n_e} 
abla_{r_i}^2 + \, V(r_1,...,r_{n_e},R_1,...,R_{n_N}) \ \mathcal{V} &= \, q_e^2 \left[ - \, \sum_{i}^{n_N} \sum_{j}^{n_e} rac{\mathcal{Z}_i}{|R_i - r_j|} + \, \sum_{i,j > i}^{n_N} rac{\mathcal{Z}_i \mathcal{Z}_j}{|R_i - R_j|} + \, \sum_{i,j > i}^{n_e} rac{1}{|r_i - r_j|} 
ight] \, . \end{aligned}$$

Nuclei are heavier so slower

$$egin{aligned} T_{N} &= -rac{\hbar^{2}}{2}\sum_{i}^{n_{N}}rac{1}{m_{i}}
abla^{2}_{R_{i}} ext{ neglected} 
ightarrow \Psi_{lpha}(r,R) &pprox arphi_{lphalpha}(R)\phi_{lpha}(r;R) \ &H_{e}\phi_{lpha}(r;R) = V(R)\phi_{lpha}(r;R) \end{aligned}$$

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#### Nuclei are heavier so slower

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### Failing of Born-Oppenheimer approx.

Before considering Born-Oppenheimer approximation

$$egin{aligned} \Psi_eta(r,R) &= \sum_lpha arphi_{etalpha}(R) \phi_lpha(r;R) ext{ with } arphi_{etalpha} ext{ satisfying} \ &[\langle \phi_lpha(r;R) | T_N | \phi_{lpha'}(r;R) 
angle + V(R) - E] \,arphi_{etalpha}(R) = \ &- \sum_{lpha 
eq lpha'} \langle \phi_lpha(r;R) | T_N | \phi_{lpha'}(r;R) 
angle arphi_{etalpha'}(R) \end{aligned}$$



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angle arphi_{etalpha'}(R) \end{aligned}$$

To evaluate kinetic couplings we use commutator  $[T_N, H_e]$ 

 $\langle \phi_{\alpha} | [T_{N}, H_{e}] | \phi_{\alpha'} \rangle = (V_{\alpha'} - V_{\alpha}) \langle \phi_{\alpha} | T_{N} | \phi_{\alpha'} \rangle + \langle \phi_{\alpha} | \phi_{\alpha'} \rangle T_{N} V_{\alpha}$ 

$$\langle \phi_{lpha} | \phi_{lpha'} 
angle = \mathbf{0} 
ightarrow \langle \phi_{lpha} | T_{N} | \phi_{lpha'} 
angle = rac{\langle \phi_{lpha} | [T_{N}, H_{e}] | \phi_{lpha'} 
angle}{(V_{lpha'} - V_{lpha})}$$

## Failing of Born-Oppenheimer approx. Before considering Born-Oppenheimer approximation $\Psi_{\beta}(r, R) = \sum \varphi_{\beta\alpha}(R) \phi_{\alpha}(r; R)$ with $\varphi_{\beta\alpha}$ satisfying $[\langle \phi_{\alpha}(\mathbf{r}; \mathbf{R}) | T_{N} | \phi_{\alpha'}(\mathbf{r}; \mathbf{R}) \rangle + V(\mathbf{R}) - \mathbf{E}] \varphi_{\beta\alpha}(\mathbf{R}) =$ $1-\sum\langle \phi_lpha(\pmb{r};\pmb{R})|T_{\pmb{N}}|\phi_{lpha'}(\pmb{r};\pmb{R}) angle arphi_{etalpha'}(\pmb{R})$ $\alpha \neq \alpha'$ To evaluate kinetic couplings we use commutator $[T_N, H_e]$ $\langle \phi_{\alpha} | [T_{N}, H_{e}] | \phi_{\alpha'} \rangle = (V_{\alpha'} - V_{\alpha}) \langle \phi_{\alpha} | T_{N} | \phi_{\alpha'} \rangle + \langle \phi_{\alpha} | \phi_{\alpha'} \rangle T_{N} V_{\alpha'}$ $\langle \phi_{\alpha} | \phi_{\alpha'} \rangle = \mathbf{0} \rightarrow \langle \phi_{\alpha} | T_{N} | \phi_{\alpha'} \rangle = \frac{\langle \phi_{\alpha} | [T_{N}, H_{e}] | \phi_{\alpha'} \rangle}{(V_{\alpha'} - V_{a})}$

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Main steps for the calculation of potential interaction

• Electronic structure package - Molpro, Molcas, Gaussian



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#### In our studies

- (Mg, Sr, H Dunning basis, Ca Ahlrichs "def2" basis, Ba - Pseudopotential (10 e<sup>-</sup>) + basis
- Configuration Interaction (CI) calculations

### Diabatization with MOLPRO



#### Method

Quasi-diabatization method based on the maximization of the overlap between states at a certain geometry with respect to the overlap at a reference geometry

#### Cons

- The method cannot be manipulate
- Presence of ionic crossing state troubles the method
- Diabatic states are not smooth
- Expensive method 2 steps:
  - Maximization of overlap
  - DDR to find transformation Double number of points

Pros

Method is automatic



#### Method

NACME are obtained as the solution of a second order diferential equation which only requires the overlap matrix and its derivatives

#### Cons

- Overlap matrix requires manipulation
- Presence of ionic crossing state troubles the method
- It requires the computation of many points

#### Pros

#### Diabatic states are smooth

### **Diabatic states**





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### **Diabatic states - MOLPRO**





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#### My collaborators:

Octavio Roncero Villa - CSIC Alfredo Aguado Gómez - UAM

#### Finantial support:

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